

# The Schrödinger Equation

Originally appeared at:

<http://behindtheguesses.blogspot.com/2009/05/schrodinger-equation.html>

Eli Lansey — [elansy@gmail.com](mailto:elansy@gmail.com)

May 26, 2009



---

**Update:** A corrected and improved version of this post is now up at:  
<http://behindtheguesses.blogspot.com/2009/06/schrodinger-equation-corrections.html>

notElon asked me to discuss, and to try and derive the Schrödinger equation, so I'll give it a shot. This derivation is partially based on Sakurai,[1] with some differences.

## A brief walk through classical mechanics

Say we have a function of  $f(x)$  and we want to translate it in space to a point  $(x+a)$ . To do this, we'll find a "space translation" operator  $\mathcal{S}_a$  which, when applied to  $f(x)$ , gives  $f(x+a)$ . That is,

$$f(x+a) = \mathcal{S}_a f(x) \tag{1}$$

We'll expand  $f(x+a)$  in a Taylor series:

$$\begin{aligned} f(x+a) &= f(x) + a \frac{df(x)}{dx} + \frac{a^2}{2!} \frac{d^2 f(x)}{dx^2} + \dots \\ &= \left[ 1 + a \frac{d}{dx} + \frac{a^2}{2!} \frac{d^2}{dx^2} + \dots \right] f(x) \end{aligned} \tag{2}$$

which can be simplified using the series expansion of the exponential<sup>1</sup> to

$$e^{[a \frac{d}{dx}]} f(x) \tag{3}$$

from which we can conclude that

$$\mathcal{S}_a = e^{[a \frac{d}{dx}]} \tag{4}$$

If you do a similar thing with rotations around the  $z$ -axis, you'll find that the rotation operator is

$$\mathcal{R}_\theta = e^{\theta L_z}, \tag{5}$$

where  $L_z$  is the  $z$ -component of the angular momentum.

Comparing (4) and (5), we see that both have an exponential with a parameter (distance or angle) multiplied by something ( $\frac{d}{dx}$  or  $L$ ). We'll call the something the "generator of the transformation." So, the generator of space translation is  $\frac{d}{dx}$  and the generator of rotation is  $L$ . So, we'll write an arbitrary transformation operator  $\mathcal{O}$  through a parameter  $\alpha$  as

$$\mathcal{O}_\alpha = e^{\alpha G} \tag{6}$$

where  $G$  is the generator of this particular transformation.<sup>2</sup> See [2] for an example with Lorentz transformations.

---

<sup>1</sup> $e^x = \sum_{n=0}^{\infty} \frac{x^n}{n!} = 1 + x + \frac{x^2}{2!} + \dots$

<sup>2</sup>There are other ways to do this, differing by factors of  $i$  in the definition of the generators and in the construction of the exponential, but I'm sticking with this one for now.

## From classical to quantum

In classical dynamics, the time derivative of a quantity  $f$  is given by the Poisson bracket:

$$\frac{df}{dt} = \{f, H\} \quad (7)$$

where  $H$  is the classical Hamiltonian of the system and  $\{ , \}$  is shorthand for a messy equation.[3] In quantum mechanics this equation is replaced with

$$\frac{df}{dt} = i\hbar[f, \mathcal{H}] \quad (8)$$

where the square brackets signify a commutation relation and  $\mathcal{H}$  is the quantum mechanical Hamiltonian.[4] This holds true for any quantity  $f$ , and  $i\hbar$  is a number which commutes with everything, so we can argue that the quantum mechanical Hamiltonian operator is related to the classical Hamiltonian by

$$H = i\hbar\mathcal{H} \Rightarrow \mathcal{H} = -iH/\hbar \quad (9)$$

specifically.

Additionally, we can extend from here that any quantum operator  $\mathcal{G}$  is written in terms of its classical counterpart  $G$  by

$$\mathcal{G} = -iG/\hbar. \quad (10)$$

So, using (4) the quantum mechanical space translation operator is given by

$$\mathcal{S}_a = e^{[-i\frac{a}{\hbar}\frac{d}{dx}]} \quad (11)$$

and, using (5), the rotation operator by

$$\mathcal{R}_\theta = e^{-i\frac{\theta}{\hbar}L_z} \quad (12)$$

or, from (6) any arbitrary (unitary) transformation,  $\mathcal{U}$ , can be written as

$$\mathcal{U} = e^{-i\frac{\alpha}{\hbar}G}, \quad (13)$$

where  $G$  is (an Hermitian operator and is) the classical generator of the transformation.

## Time translation of a quantum state

Consider a quantum state at time  $t$  described by the wavefunction  $\psi(\vec{r}, t)$ . To see how the state changes with time, we want to find a “time-translation” operator  $\mathcal{T}_{\Delta t}$  which, when applied to the state  $\psi(\vec{r}, t)$ , will give  $\psi(\vec{r}, t + \Delta t)$ . That is,

$$\psi(\vec{r}, t + \Delta t) = \mathcal{T}_{\Delta t}\psi(\vec{r}, t). \quad (14)$$

From our previous discussion we know that if we know the classical generator of time translation we can write  $\mathcal{T}$  using (13). Well, classically, the generator of time translations is the Hamiltonian![5] So we can write

$$\mathcal{T}_{\Delta t} = e^{-i\frac{\Delta t}{\hbar}H} \quad (15)$$

and (14) becomes

$$\psi(\vec{r}, t + \Delta t) = e^{-i\frac{\Delta t}{\hbar}H} \psi(\vec{r}, t). \quad (16)$$

This holds true for any time translation, so we'll consider a small time translation and expand (16) using a Taylor expansion<sup>3</sup> dropping all quadratic and higher terms:

$$\psi(\vec{r}, t + \Delta t) \approx \left[ 1 - i\frac{\Delta t}{\hbar}H + \dots \right] \psi(\vec{r}, t) \quad (17)$$

Moving things around gives

$$H\psi(\vec{r}, t) = i\hbar \left[ \frac{\psi(\vec{r}, t + \Delta t) - \psi(\vec{r}, t)}{\Delta t} \right] \quad (18)$$

In the limit  $\Delta t \rightarrow 0$  the righthand side becomes a partial derivative giving the Schrödinger equation

$$H\psi(\vec{r}, t) = i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t} \quad (19)$$

For a system with conserved total energy, the classical Hamiltonian is the total energy

$$H = \frac{\vec{p}^2}{2m} + V \quad (20)$$

which, making the substitution for quantum mechanical momentum  $\vec{p} = i\hbar\nabla$  and substituting into (19) gives the familiar differential equation form of the Schrödinger equation

$$-\frac{\hbar^2}{2m}\nabla^2\psi(\vec{r}, t) + V\psi(\vec{r}, t) = i\hbar\frac{\partial \psi(\vec{r}, t)}{\partial t} \quad (21)$$

## References

- [1] J.J. Sakurai. *Modern Quantum Mechanics*. Addison-Wesley, San Francisco, CA, revised edition, 1993.
- [2] J.D. Jackson. *Classical Electrodynamics*. John Wiley & Sons, Inc., 3rd edition, 1998.
- [3] L.D. Landau and E.M. Lifshitz. *Mechanics*. Pergamon Press, Oxford, UK, 3rd edition, 1976.
- [4] L.D. Landau and E.M. Lifshitz. *Quantum Mechanics*. Butterworth-Heinemann, Oxford, UK, 3rd edition, 1977.
- [5] H. Goldstein, C. Poole, and J. Safko. *Classical Mechanics*. Cambridge University Press, San Francisco, CA, 3rd edition, 2002.

---

<sup>3</sup>Kind of the reverse of how we got to this whole exponential notation in the first place...